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# VELOCITY DISTRIBUTION OF METEOROIDS COLLIDING WITH PLANETS AND SATELLITES. II. NUMERICAL RESULTS 

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In the first part of the paper we proposed algorithm for describing velocity distribution of meteoroids colliding with planets and satellites. In the present part we show numerical characteristics of the distribution function. Namely, for each of terrestrial planets and their satellites we consider a swarm of encountering particles of asteroidal origin. They form a field of relative collisional velocities $v$. We consider momenta $\mu_{k}$ (mathematical expectation of $v^{k}$ ), $k=-1,1,2,3,4$. The data are calculated under two different assumptions: taking into account gravitation of target body or without it. The main results are presented in a series of tables each containing five numbers and several useful functions of them.

KEY WORDS Meteoroids, velocity distribution, collision

## 1 INTRODUCTION

Collisions in the Solar System play an important role in its history. Consequencies of such events depend essentially on relative velocity of the impactors. In Part I of the present paper (Kholshevnikov and Shor, 1994) we described the method for obtaining the velocity distribution function of meteoroids colliding with planets and satellites. The main idea of the method is as follows. Let us fix a body-target, say the $s$-th major planet $Q_{z}$, and a set of potential projectiles, say minor planets $Q$. Choose among all the numbered minor planets (Batrakov, 1992) those which have the semi-major axis, $a$, and eccentricity, $e$, satisfying the inequalities

$$
\begin{equation*}
a(1-e)<a_{s}\left(1+e_{s}\right), \quad a(1+e)>a_{s}\left(1-e_{s}\right) \tag{1}
\end{equation*}
$$

where the elements with index $s$ refer to $Q_{s}$.
As the lines of nodes and apsides rotate, the orbits of $Q$ and $Q_{s}$ intersect each other from time to time. The relative velocities at intersection points can be calculated without difficulties. When averaging over all possible intersection points and

Table 1. Velocity distribution characteristics when colliding with Mercury; gravitation is not taken into account; the last row represents $\lambda_{k}$ for Maxwell's distribution

| $k$ | -1 | 1 | 2 | $\rho$ | 4 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\mu_{k}$ | 0.03588 | 33.08 | 1281 | 56090 | 2684000 |
| $\nu_{k}$ | 27.87 | 33.08 | 35.79 | 38.28 | 40.48 |
| $\lambda_{k}$ | 0.8426 | 1 | 1.082 | 1.157 | 1.224 |
| $\lambda_{k}$ | 0.7854 | 1 | 1.085 | 1.162 | 1.233 |
|  |  | $N=6, \sigma=13.66 \mathrm{~km} / \mathrm{s}$ |  |  |  |
|  |  |  |  |  |  |

all selected asteroids, we obtain a set of $\mu_{k}=$ mathematical expectation of $v^{k}, v$ being relative collisional velocity. According to the Carleman theorem (see, e.g., Prokhorov and Rosanov, 1973, §4.3), the distribution function is uniquely determined by its momenta $\mu_{k}$. We expect that $\mu_{k}$ for meteoroids of asteroidal nature differ very little from those for asteroids themselves.

More detailed discussion and the description of algorithms one can find in Part I of the present paper. The closed formulae for $v$ and numerical process for averaging over all intersection points are also given there. For $\mu_{2}, \mu_{4}$ the averaging can be fulfilled also in the analytical form. The results coincide at least to four decimals.

## 2 MERCURY

For $s=1$ (Mercury) there are 6 astcroids catalogued in Batrakov (1992) satisfying the inequalities (1).

Table 1 contains five momenta $\mu_{k}=$ mathematical expectation of $v^{k}$ (first row), $v$ being planetocentric velocity of a collider calculated without taking into account Mercurian gravity. In the second and the third rows, correspondingly, $\nu_{k}=\left(\mu_{k}\right)^{1 / k}$ and $\lambda_{k}=\nu_{k} / \nu_{1}$ are given. The last row represents $\lambda_{k}$ for Maxwell's velocity distribution designated $\tilde{\lambda_{k}}$ (comparing $\lambda_{k}$ and $\tilde{\lambda_{k}}$ one can notice their proximity).

In this paper distances are measured in km , velocities - in $\mathrm{km} / \mathrm{s}$. So the first line contains the quantities of different dimensions $(\mathrm{km} / \mathrm{s})^{k}$; the third and fourth line quantities are dimensionless.

Below Table 1 we give $N=$ number of minor planets - potential colliders and $\sigma=\sqrt{\mu_{2}-\mu_{1}^{2}}=$ the mean squared deviation of velocity.

We described the collision velocity distribution provided that the planet's gravity is negligible. In reality it is a planetocentric velocity distribution in the meteoroid swarm at the border of the planet's sphere of action. Table 1a contains the same quantities as Table 1 (except Maxwell's momenta) calculated with due regard of planet's attraction. More precisely Table la contains characteristics of planetocentric velocity field at a distance $R$ ( $R$ being the radius of the planet) from the centre of Mercury. The corresponding correction depends on the parabolic velocity $v_{0}(R)$ only - see formula (12) in Kholshevnikov and Shor (1994). When calculating, the rotation of the planet was not taken into account. The corresponding correction

Table 1a. Velocity distribution characteristics when colliding with Mercury; gravitation is taken into account

| $k$ | -1 | 1 | 2 | 9 | 4 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\mu_{k}$ | 0.03527 | 33.40 | 1299 | 56990 | 2731000 |
| $\nu_{k}$ | 28.35 | 33.40 | 36.04 | 38.48 | 40.65 |
| $\lambda_{k}$ | 0.8489 | 1 | 1.079 | 1.152 | 1.217 |
|  | $N=6, R=2439.7 \mathrm{~km}, v_{0}$ | $=4.249 \mathrm{~km} / \mathrm{s}, \sigma=13.54 \mathrm{~km} / \mathrm{s}$ |  |  |  |

Table 1b. Minimal and maximal velocities and length of interval of possible collision for Mercury-crossers; the second and the third columns correspond to the case when gravitation is not taken into account; the fourth and the fifth columns, when gravitation is taken into account

| $N^{*}$ | $v_{\min }$ | $v_{\max }$ | $v_{\min }^{*}$ | $v_{\max }^{*}$ | $\Omega_{2}-\Omega_{1}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 1566 | 34.99 | 54.31 | 35.25 | 54.48 | $180^{\circ}$. |
| 2101 | 19.03 | 23.25 | 19.50 | 23.63 | 38.39 |
| 2212 | 17.31 | 34.19 | 17.82 | 34.45 | 97.49 |
| 2340 | 16.86 | 17.21 | 17.38 | 17.73 | 11.56 |
| 3200 | 43.65 | 64.12 | 43.85 | 64.26 | 180. |
| 3838 | 34.35 | 35.96 | 34.62 | 36.21 | 32.24 |
|  |  | $\Delta v=\left(v_{\max }-v_{\min }\right)_{\operatorname{mean}}=10.48 \mathrm{~km} / \mathrm{s}$ |  |  |  |
|  | $\Delta v^{*}=\left(v_{\max }^{*}-v_{\min }^{*}\right)_{\operatorname{mean}}=10.39 \mathrm{~km} / \mathrm{s}$ |  |  |  |  |
|  |  |  |  |  |  |

depends on the velocity direction and place of impact, but it is rather small even for the Earth and Mars.

Now we return to the mean squared deviation $\sigma$. Statistical rules demand a factor $b=\sqrt{N /(N-1)}$ when estimating $\sigma$ from a sample containing $N$ values of a random variable. Since $N=6$, the factor $b=1.095$ is important. But we ought to remember that each selected asteroid brings a set of $20-360$ points. So in reality we deal with a quantity $N^{\prime}=1076$ and the discussed factor may be omitted.

The above numbers $20-360$ and 1076 need explanation. If the elements of $Q$ satisfy inequalities more restrictive in comparison with (1),

$$
\begin{equation*}
a(1-e)<a_{s}\left(1-e_{s}\right), \quad a(1+e)>a_{s}\left(1+e_{s}\right) \tag{2}
\end{equation*}
$$

then a certain position of apsidal line, leading to intersection of the orbits, corresponds to each value of $\Omega$ (the longitude of ascending node). We calculated values of $v$ for a set of $\Omega=0^{\circ}, 1^{\circ}, \ldots, 359^{\circ}$. If the inequalities (1), but not (2) are satisfied, then not all values of $\Omega$ lead to intersection. In this case $\Omega \in\left[\Omega_{1}, \Omega_{2}\right] \cup\left[-\Omega_{2},-\Omega_{1}\right]$ and $0^{\circ} \leq \Omega_{1}<\Omega_{2} \leq 180^{\circ}$. If $\Omega_{2}-\Omega_{1} \geq 10^{\circ}$, one chooses a $1^{\circ}$ step in $\Omega$. If not, one takes exactly 10 steps.

In Table 1b we give the number $N^{*}$ of the selected minor planet (the first column); minimal and maximal velocities over all possible intersection points when

Table 2. Velocity distribution characteristics when colliding with Venus; gravitation is not taken into account

| $k$ | -1 | 1 | 2 | $s$ | 4 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\mu_{k}$ | 0.05448 | 21.02 | 501.5 | 13390 | 393000 |
| $\nu_{k}$ | 18.36 | 21.02 | 22.39 | 23.75 | 25.04 |
| $\lambda_{k}$ | 0.8734 | 1 | 1.065 | 1.130 | 1.191 |
|  |  | $N=25, \sigma=7.73 \mathrm{~km} / \mathrm{s}$ |  |  |  |

Table 2a. Velocity distribution characteristics when colliding with Venus; gravitation is taken into account

| $k$ | -1 | 1 | 2 | 9 | 4 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\mu_{k}$ | 0.04566 | 23.68 | 608.8 | 17000 | 512200 |
| $\nu_{k}$ | 21.90 | 23.68 | 24.67 | 25.71 | 26.75 |
| $\lambda_{k}$ | 0.9250 | 1 | 1.042 | 1.086 | 1.130 |
|  | $N=25, R=6051.8 \mathrm{~km}, \nu_{0}=10.362 \mathrm{~km} / \mathrm{s}, \sigma=6.95 \mathrm{~km} / \mathrm{s}$ |  |  |  |  |

Table 3. Velocity distribution characteristics when colliding with the Earth; gravitation is not taken into account

| $k$ | -1 | 1 | 2 | 3 | 4 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\mu_{k}$ | 0.06489 | 18.15 | 379.4 | 8866 | 225000 |
| $\nu_{k}$ | 15.41 | 18.15 | 19.48 | 20.70 | 21.78 |
| $\lambda_{k}$ | 0.8490 | 1 | 1.073 | 1.140 | 1.200 |
|  |  | $N=45, \sigma=7.07 \mathrm{~km} / \mathrm{s}$ |  |  |  |

Table 3a. Velocity distribution characteristics when colliding with the Earth; gravitation is taken into account

| $k$ | -1 | 1 | 2 | 9 | 4 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\mu_{k}$ | 0.04964 | 21.65 | 504.6 | 12620 | 335600 |
| $\nu_{k}$ | 20.15 | 21.65 | 22.46 | 23.28 | 24.07 |
| $\lambda_{k}$ | 0.9306 | 1 | 1.038 | 1.075 | 1.112 |
|  | $N=45, R=6371.0 \mathrm{~km}, \nu_{0}=11.186 \mathrm{~km} / \mathrm{s}, \sigma=5.99 \mathrm{~km} / \mathrm{s}$ |  |  |  |  |
|  |  |  |  |  |  |

Table 3b. Minimal and maximal velocities and length of interval of possible collision for Earth-crossers; the second and the third columns correspond to the case when gravitation is not taken into account; the fourth and the fifth columns, when gravitation is taken into account

| $N^{*}$ | $v_{\text {min }}$ | $v_{\max }$ | $v_{\text {min }}^{*}$ | $v_{\text {max }}^{*}$ | $\Omega_{2}-\Omega_{1}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 1566 | 29.10 | 30.36 | 31.18 | 32.35 | $180^{\circ}$ |
| 1620 | 11.19 | 11.97 | 15.82 | 16.39 | 180 |
| 1685 | 12.59 | 13.52 | 16.84 | 17.55 | 180 |
| 1862 | 16.43 | 17.41 | 19.88 | 20.69 | 180 |
| 1863 | 15.57 | 16.40 | 19.17 | 19.85 | 180 |
| 1864 | 21.64 | 22.55 | 24.36 | 25.17 | 180 |
| 1865 | 16.06 | 17.02 | 19.57 | 20.36 | 180 |
| 1866 | 25.14 | 25.65 | 27.52 | 27.99 | 180 |
| 1981 | 27.56 | 28.38 | 29.74 | 30.51 | 180 |
| 2062 | 10.69 | 11.35 | 15.47 | 15.94 | 180 |
| 2063 | 11.19 | 12.11 | 15.82 | 16.49 | 180 |
| 2100 | 13.34 | 14.51 | 17.41 | 18.32 | 180 |
| 2101 | 24.43 | 25.45 | 26.87 | 27.80 | 180 |
| 2102 | 33.16 | 34.10 | 34.99 | 35.89 | 180 |
| 2135 | 17.60 | 18.30 | 20.86 | 21.45 | 180 |
| 2201 | 19.88 | 20.87 | 22.81 | 23.68 | 180 |
| 2212 | 28.45 | 29.49 | 30.57 | 31.54 | 180 |
| 2329 | 19.87 | 20.58 | 22.80 | 23.43 | 180 |
| 2340 | 12.12 | 13.35 | 16.49 | 17.41 | 180 |
| 3103 | 13.64 | 14.18 | 17.64 | 18.06 | 180 |
| 3200 | 32.82 | 34.09 | 34.68 | 35.88 | 180 |
| 3360 | 23.42 | 24.29 | 25.95 | 26.74 | 180 |
| 3361 | 8.60 | 9.61 | 14.11 | 14.75 | 180 |
| 3362 | 14.61 | 15.66 | 18.40 | 19.25 | 180 |
| 3554 | 14.02 | 14.80 | 17.94 | 18.55 | 180 |
| 3671 | 10.69 | 11.16 | 15.47 | 15.80 | 78.62 |
| 3752 | 29.58 | 30.24 | 31.62 | 32.24 | 145.95 |
| 3753 | 18.00 | 19.03 | 21.19 | 22.08 | 180 |
| 3757 | 6.60 | 6.62 | 12.99 | 13.00 | 12.91 |
| 3838 | 26.33 | 27.32 | 28.61 | 29.52 | 180 |
| 4015 | 8.28 | 9.38 | 13.92 | 14.60 | 100.81 |
| 4034 | 14.30 | 15.28 | 18.15 | 18.94 | 180 |
| 4179 | 11.51 | 12.86 | 16.05 | 17.04 | 180 |
| 4183 | 16.89 | 17.89 | 20.26 | 21.10 | 180 |
| 4197 | 24.15 | 25.12 | 26.62 | 27.50 | 180 |
| 4257 | 24.19 | 24.71 | 26.65 | 27.12 | 180 |
| 4341 | 20.75 | 21.70 | 23.57 | 24.41 | 180 |
| 4450 | 17.63 | 18.62 | 20.88 | 21.72 | 180 |
| 4486 | 16.49 | 17.54 | 19.93 | 20.80 | 180 |
| 4544 | 10.10 | 10.85 | 15.07 | 15.59 | 180 |
| 4581 | 10.64 | 11.65 | 15.44 | 16.15 | 180 |
| 4660 | 6.07 | 7.51 | 12.73 | 13.47 | 180 |
| 4769 | 15.13 | 16.15 | 18.81 | 19.64 | 180 |
| 4953 | 23.29 | 24.19 | 25.84 | 26.65 | 180 |
| 5011 | 12.54 | 13.55 | 16.80 | 17.57 | 180 |
| $\begin{aligned} \Delta v & =\left(v_{\text {max }}-v_{\text {min }}\right)_{\text {mean }}=0.91 \mathrm{~km} / \mathrm{s} \\ \Delta v^{*} & =\left(v_{\text {max }}^{*}-v_{\text {min }}^{*}\right)_{\text {mean }}=0.74 \mathrm{~km} / \mathrm{s} \end{aligned}$ |  |  |  |  |  |

Table 4. Velocity distribution characteristics when colliding with Mars; gravitation is not taken into account

| $k$ | -1 | 1 | 2 | 9 | 4 |
| :--- | :--- | :--- | :---: | :---: | :---: |
| $\mu_{k}$ | 0.1002 | 11.87 | 164.4 | 2595 | 45590 |
| $\nu_{k}$ | 9.983 | 11.87 | 12.82 | 13.74 | 14.61 |
| $\lambda_{k}$ | 0.8411 | 1 | 1.080 | 1.158 | 1.231 |
|  |  | $N=165, \sigma=4.85 \mathrm{~km} / \mathrm{s}$ |  |  |  |

Table 4a. Velocity distribution characteristics when colliding with Mars; gravitation is taken into account

| $k$ | -1 | 1 | 2 | $\rho$ | 4 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\mu_{k}$ | 0.08551 | 13.03 | 189.6 | 3067 | 54540 |
| $\nu_{k}$ | 11.69 | 13.03 | 13.77 | 14.53 | 15.28 |
| $\lambda_{k}$ | 0.8974 | 1 | 1.057 | 1.115 | 1.173 |
|  | $N=165, R=3390.0 \mathrm{~km}, \nu_{0}=5.029 \mathrm{~km} / \mathrm{s}, \sigma=4.45 \mathrm{~km} / \mathrm{s}$ |  |  |  |  |

gravity is not taken into account (the second and the third columns): these quantities obtained in consideration of planet's gravity (the fourth and the fifth columns); $\Omega_{2}-\Omega_{1}$ in degrees. In particular, we can extract the number $N^{\prime}=1076$ from the last column.

## 3 VENUS - EARTH - MARS

Velocity distribution characteristics for the $s$-th planet are listed in the Table having the number $s$ (negligible gravity; the letter "a" attached to the number indicates that the gravity is taken into account). To save room, we present the Tables of the 1b type for Mercury and the Earth only.

Calculation of the data of the Tables is described in the previous section.
In the case of Venus, the difference $v_{\max }-v_{\min }$ varies within narrow limits from 0.15 to 0.71 , and its averaged value is $\Delta v=0.46$, that is $2 \%$ from $\mu_{1}=21$. This fact reflects the following property. If $e_{s}=0$, then all positions of the line of nodes are equivalent and $v_{\max }=v_{\min }$ for each planet-crosser. For Venus $e_{2}=0.007$. In this case we may estimate $N^{\prime} \approx 4 N=100$ and neglect the factor $b \approx 1.005$.

For the Earth $v_{\max }-v_{\min }$ varies between $0.02 \mathrm{~km} / \mathrm{s}$ and $1.44 \mathrm{~km} / \mathrm{s}, \Delta v=$ $0.91 \mathrm{~km} / \mathrm{s}$.

For Mars the limits of $v_{\max }-v_{\min }$ are 0.01 and $6.84, \Delta v=2.61$.
For the Earth and especially for Mars, the eccentricities as well as the number of selected asteroids are larger than for Venus. So we put $b=1$ all the more.

Table 5. Velocity distribution characteristics when colliding with the Moon; gravitation is not taken into account

| $k$ | -1 | 1 | 2 | 9 | 4 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\mu_{k}$ | 0.06433 | 18.25 | 382.6 | 8957 | 227700 |
| $\nu_{k}$ | 15.54 | 18.25 | 19.56 | 20.77 | 21.84 |
| $\lambda_{k}$ | 0.8518 | 1 | 1.072 | 1.138 | 1.197 |
|  |  | $N=45, \sigma=7.04 \mathrm{~km} / \mathrm{s}$ |  |  |  |

Table 5a. Velocity distribution characteristics when colliding with the Moon; gravitation is taken into account

| $k$ | -1 | 1 | 2 | $s$ | 4 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\mu_{k}$ | 0,06315 | 18.43 | 388.2 | 9112 | 232000 |
| $\nu_{k}$ | 15.83 | 18.43 | 19.70 | 20.89 | 21.95 |
| $\lambda_{k}$ | 0.8593 | 1 | 1.069 | 1.133 | 1.191 |
|  | $N=45, R=1737.4 \mathrm{~km}, u_{0}$ | $=2.376 \mathrm{~km} / \mathrm{s}, \sigma=6.98 \mathrm{~km} / \mathrm{s}$ |  |  |  |

## 4 THE MOON, PHOBOS AND DEIMOS

Velocity distribution characteristics for a satellite are determined by the same set of asteroids as for the corresponding planet. For the Moon they are represented in Table 5 and 5a, for Phobos and Deimos, in Table 6 and 7. The gravity field of Martian satellites is negligible for our purpose: the parabolic velocity $u_{0}$ on their surfaces is less than $0.01 \mathrm{~km} / \mathrm{s}$. Remember that we deal with squared velocities and $u_{0}^{2} / \mu_{2}<10^{-6}$.

The coincidence of $\mu_{2}$ calculated by means of numerical and analytical averaging is as perfect as in the case of planets.

## 5 DISCUSSION

1. It should be noted that the results of our calculations of minimal and maximal collisional velocities of minor planets with Mars (to save room, we do not give them in a special Table) systematically differ from those published by Steel (1985). According to our calculations, the lower boundary of collisional velocities is, as a rule, by several hundred meters per second greater and the upper boundary by several hundred meters per second smaller. Variation of this difference lies within limits from 0 to $1.5 \mathrm{~km} / \mathrm{s}$. The minor planet (1727) Mette is a typical example. The perihelion distance of this planet is equal to $1.665 \mathrm{a} . \mathrm{u}$. , that is very close to aphelion distance of Mars ( 1.666 a.u.). Collision of a body having the same orbital elements $a, e, i$ as Mette has with Mars is possible only within small vicinity

Table 6. Velocity distribution characteristics when colliding with Phobos

| $k$ | -1 | 1 | 2 | 3 | 4 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\mu_{k}$ | 0.09315 | 12.43 | 176.9 | 2843 | 50710 |
| $\nu_{k}$ | 10.74 | 12.43 | 13.30 | 14.17 | 15.01 |
| $\lambda_{k}$ | 0.8637 | 1 | 1.070 | 1.140 | 1.207 |
|  |  | $N=165, \sigma=4.73 \mathrm{~km} / \mathrm{s}$ |  |  |  |

Table 7. Velocity distribution characteristics when colliding with Deimos

| $k$ | -1 | 1 | 2 | 9 | 4 |
| :--- | :---: | :---: | :---: | :---: | :---: |
| $\mu_{k}$ | 0.09754 | 12.06 | 168.6 | 2683 | 47550 |
| $\nu_{k}$ | 10.25 | 12.06 | 12.99 | 13.90 | 14.77 |
| $\lambda_{k}$ | 0.8501 | 1 | 1.077 | 1.152 | 1.224 |
|  |  | $N=165, \sigma=4.82 \mathrm{~km} / \mathrm{s}$ |  |  |  |

of perihelion. Our calculations give for this planet $\Omega_{2}-\Omega_{1}=6^{\circ} .00$. Moreover, collision velocity with Mars in any point of this vicinity is practically the same and equal to $9.38 \mathrm{~km} / \mathrm{s}$. In Steel's paper for this planet $v_{\min }$ and $v_{\max }$ are given as $8.7 \mathrm{~km} / \mathrm{s}$ and $10.2 \mathrm{~km} / \mathrm{s}$, correspondingly. Such a distinction appears, first of all, as a consequence of different setting up the problem. The orbit of a major planet, say, Mars is treated here as a fixed one with elements referred to the contemporary epoch. One examine impact velocities of the flux of particles having the same fixed orbital elements $a, e, i$ (inclination to the Mars orbit) as the chosen minor planet has and uniformly distributed directions of nodes and perihelions. Minimal and maximal relative velocities over all intersection points with the Mars orbit are found. Final results represent an estimation of impact velocities of meteoroids of asteroidal nature with the planet for the present epoch.

On the contrary, Steel's results correspond to the assumption that inclinations to the invariant plane conserve whereas distribution of nodes and perihelions of both Mars and the minor planet is uniform. This implies a variation of $i$ and extension of the set of velocity values. Velocity distribution turns out to be valid over a period of perihelion revolution of Mars, i.e. about one hundred of thousand years.
2. When examining the Tables above, we can extract the following conclusions.
(1) Compare the third rows of the Tables numbered $s$, sa $(s=1, \ldots, 7)$ with the fourth one of Table 1 . We see that velocity distributions are similar to Maxwell's one, but definitely do not coincide with it. We shall try to find the analytical form of distribution functions in the next paper.
(2) Compare $\nu_{k}, \sigma$ for fixed $k$ from the set of Tables s or those of sa $(s=1, \ldots, 4)$. In both cases $\nu_{k}, \sigma$ decrease with $s$. The reason is evident: the farther from the

Sun, the greater the potential and the smaller the kinetic energy of projectiles are.
(3) Compare the Table $s(s=1, \ldots, 5)$ with the Table sa. Confrontation of the second and the third rows shows that planet's or Moon's gravitation augments $\nu_{k}$ (this is obvious without calculations) and makes $\lambda_{k}$ nearer to 1. In other words, gravitation of a target body makes the random variable $v$ "less random". This is reflected also in diminishing mean squared deviation $\sigma$.
(4) Compare the quantity $\nu_{k}$ from Table $s$ with that from Table sa, designating latter as $\alpha_{k}$. Then we have (Kholshevnikov and Shor, 1994):

$$
\alpha_{2}^{2}=\nu_{2}^{2}+v_{0}^{2}
$$

For $k \neq 2$ values of $\alpha_{k}$ differ from $\left(\nu_{k}^{2}+v_{2}^{2}\right)^{1 / 2}$ by several units of the third decimal.
(5) Compare Tables 3 and 5. Let us designate $\beta_{k}$ the quantity $\nu_{k}$ from Table 5, preserving notation $\nu_{k}$ for Table 3 quantity. Values of $\beta_{k}$ are greater than $\nu_{k}$ not more than by $0.06-0.13$. According to Kholshevnikov and Shor (1994),

$$
\beta_{2}^{2}=\nu_{2}^{2}+3 u^{2}
$$

$u$ being the circular geocentric velocity of the Moon. Values of $\beta_{k}$ differ from those of $\left(\nu_{k}^{2}+3 u^{2}\right)^{1 / 2}$ by $0.01(k=1,3,4)$ or $0.03(k=-1)$. This is less than precision of our data. So in future (after discovering new Earth-crossers) we do not need to fulfill new cumbersome calculations of Table 5 data: it is sufficient to use the formula

$$
\begin{equation*}
\beta_{k}^{2} \simeq \nu_{k}^{2}+3 u^{2} \tag{3}
\end{equation*}
$$

Further, using the formula

$$
\beta_{k}^{2} \simeq \nu_{k}^{2}+3 u^{2}+u_{0}^{2}
$$

for Table 5a quantity $\nu_{k}$ leads to an error in $\beta_{k}$ like $0.04(k=1,3,4)$ and 0.14 ( $k=-1$ ). Remember that $u_{0}$ is the parabolic velocity on the Moon's surface.
(6) Compare Tables 4 and 6, 7. For Phobos using the formula (3) leads to an error in $\beta_{k}$ like $0.04(k=1,3,4)$ and $0.21(k=-1)$; for Deimos, 0.02 and 0.06 .

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