Lecture 3. Helium burning (3aprocess). Degeneracy of matter in stellar interiors. Evolution of stars with M<8 M_☉. Red giants. AGB stars. Formation of planetary nebulae.

Off the main sequence.Core He burning

- After core hydrogen exhaustion: core contracts, central density and temperature rise. Star goes off the main sequence toward red giant branch as conditions for H burning in shell layers appear
- helium burning (~10% of t_{H}): starts at T~2-3x10⁸K, ρ ~10³-10⁴g/ccm
- No stable nuclei with A=5 and A=8 → first stable reaction 3 ⁴He→ ¹²C

I. ⁴He + ⁴He \leftrightarrow ⁸Be - 92 keV II. ⁸Be + ⁴He \leftrightarrow ¹²C^{*} - 0.29 MeV III. 1/2500 cases: ¹²C^{*} → ¹²C + 7.65 MeV (e⁺e⁻)

- Problem: (3m_α-m_{C12})c²=7.28 MeV. If ¹²C^{*} has no appropriate energy level, the reaction proceeds extremely slow and no other element after Helium would form in the Universe :(((.Fortunately, it has (predicted by F. Hoyle and found by W. Fowler) → resonance exists!
- Energy generation per gram $\epsilon \sim \rho^2$ (3 particles must meet)
- Immediately followed by: ⁴He+ ¹²C → ¹⁶O + 7.16 MeV Result: M<10 M_☉ - (degenerate) Carbon cores with admixture ¹⁶O 10<M<30 M_☉ - CO-cores M>30 M_☉ - Oxygen cores

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Off the main sequence: qualitative



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Off the main sequence: more quantitative

Central density increases → degeneracy of matter becomes important (especially in low mass stars as on the main sequence ρ~1/M²). Indeed, from virial theorem for Γ=5/3 Q=3/2 R<T> M = - 1/2 U ~ GM^{5/3}ρ^{1/3}, E=Q+U ~ -GM^{5/3}ρ^{1/3}

Star emits radiation so $\Delta E < 0 \rightarrow$ density must increase

- Principal consequence of increasing density Coulomb corrections + degeneracy effects to Maxwell-Boltzmann ideal gas EOS. It is the physical reason for divergence of evolution of stellar cores with different mass after the main sequence.
- If $M_{core} < M_{Ch} \sim 1.4 M_{\odot}$ (i.e. M<8-10 M_{\odot}), degeneracy wins;
- If M_{core} > M_{Ch} (i.e. M>8-10 M_☉), nuclear evolution proceeds under non-degenerate conditions up to the formation of mostly bound iron group nuclei

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Intermezzo 1: Degenerate Fermi gases



Pauli exclusion principle (1925): 2 electrons in atom cannot have identical quantum numbers – explains Periodic Table of elements! More generally: for any 2 fermions (half-integer spin: e, n, v...). Follows from (i) particles identity in quantum mechanics and (ii) odd parity of wave functions of fermions

W.Pauli (1900-1958) Degeneracy of fermions in matter at high densities or low temperatures:

Only 2 electrons with opposite spins can occupy one phase space cell

$$2\frac{d^{3}pd^{3}x}{(2\pi\hbar)^{3}} = \frac{dN}{dVdp}, \quad \int_{0}^{p_{F}} \frac{dN}{dVdp}dp = n = \frac{8\pi}{3}\frac{p_{F}^{3}}{(2\pi\hbar)^{3}}, \quad p_{F} = \pi\hbar\left(\frac{3n}{\pi}\right)^{1/3}$$
$$\varepsilon_{F} = \sqrt{p_{F}^{2}c^{2} + m^{2}c^{4}}$$

Degeneracy in ideal gas: $\varepsilon \sim kT \leq \varepsilon_F$, $T_{dgnr} \sim 1eV \times \left(\frac{\rho}{10^3 (g/cm^3)}\right)^{2/3}$ **17/10/2005** Lecture 3 $-\frac{\pi}{L}\frac{\partial}{\partial t} = \frac{P^{L}}{2m} - \frac{Ze^{L}}{L}$

E.Fermi (1901-1954)

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Intermezzo 2: Chandrasekhar limit (1930) (Nobel Prize in physics 1983, with W.Fowler)

 $\Delta x \Delta p \ge \hbar$ (*Heisenberg uncert.*, 1927)

for electrons: $p_{\lim} = p_F \sim \hbar / \Delta x = \hbar n_e^{1/3}$,

$$n_e = \frac{\rho}{m_p} Y_e, Y_e \equiv \frac{n_e}{n_b} = (1 \text{ for } H, 0.5 \text{ for } He, \frac{26}{56} \text{ for } Fe)$$

Energy:
$$\varepsilon_e \sim \varepsilon_F \simeq \frac{p_F^2}{2m_e} (non - rel.)... \simeq p_F c (ultrarel.)$$

Non-rel. pressure :
$$P_e \sim n_e \varepsilon_e = \frac{\hbar^2}{2m_e} n_e^{5/3} \propto \rho^{5/3}$$

Ultrarelativistic gas:

$$P_e \sim p_F cn_e \approx \hbar cn_e^{4/3} = K_{rel} \rho^{4/3}, \ K_{rel} = \frac{\hbar c}{m_p^{4/3}} Y_e^{4/3}$$

S.Chandrasekhar (1910-1995)

Chandrasekhar mass:

$$M_{Ch} = \left(\frac{K_{rel}}{G}\right)^{3/2} = \left(\frac{\hbar c}{m_p^{4/3}} \frac{Y_e^{4/3}}{G}\right)^{3/2} = N_{Ch} m_p Y_e^2 \approx 5.85 M_{\odot} Y_e^2 \approx 1.4 M_{\odot} (He, CO...)$$

$$N_{Ch} = \left(\frac{m_{Pl}}{m_p}\right)^3 \approx 10^{57}, \ m_{Pl} = \sqrt{\frac{\hbar c}{G}} \approx 10^{19} GeV - Plankian \ mass$$

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Why evolution of massive stars is different from evolution of low mass star and which mass determines the boundary

EOS: ideal + Fermi non – rel. $P \approx \rho \Re T + K_{nr} \rho^{5/3} \mathsf{T}$ Hydrostatic eq.: $P_c / \rho_c \sim GM / R = GM^{2/3} \rho_c^{1/3}$ Central temperature: $\Re T_c = GM^{2/3} \rho_c^{1/3} - K_{nr} \rho_c^{2/3}$ \Rightarrow at high densities T_c starts to decrease and $T \rightarrow 0$ (white dwarfs!) as $\Gamma \rightarrow 4/3$ for any $M \leq M_{Ch}$.

Fermi ultrarelativistic : $P \approx \rho \Re T + K_{rel} \rho^{4/3}$ $\Re T_c = GM^{2/3} \rho_c^{1/3} - K_{rel} \rho_c^{1/3} \Rightarrow$ if $M > M_{Ch} = (K_{rel} / G)^{3/2}$ temperature always increases $T_c \propto \rho_c^{1/3}$ (case for massive stars)







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Physical diagram of stellar evolution: ρ-T plane



Tracks in the ρ -T plane traced out by the centers of stars of various masses

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Astronomical diagram for stellar evolution: HR diagram



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Red giants





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HR diagram for stellar clusters



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Hydrogen burning time ~1/M² → Turn-off point as a measure of the cluster age (Sandage)



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AGB stars and planetary nebula formation



Thermal instability of H burning shell leads to pulsations (Mira!) and subsequent planetary nebula formation

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Zoo of planetary nebulae













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